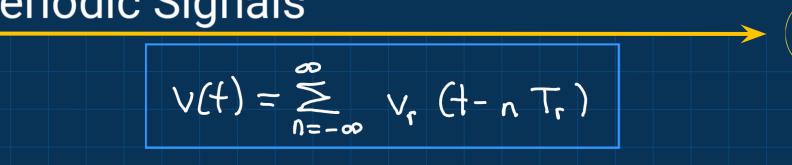
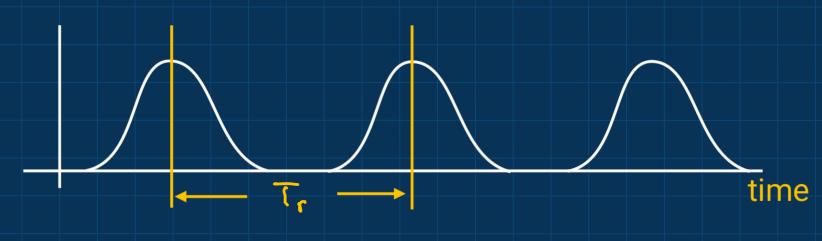


Periodic Signals







Periodic Signals



Since v(t) is repetitive we can represent v(t) as a sum of orthogonal repetitive functions. We will choose sine waves

is periodic at 📜 given m is an integer. Let

$$\omega_r = \frac{2\pi}{T_r}$$

Fourier Series



Then we can write the signal as

where C_m is complex. Multiply both sides by $e^{-jk\omega_n t}$ and integrate over a period.

$$\int_{-\tau_{1}}^{\tau_{2}} v(t) e^{-ik\omega t} dt = \sum_{m=0}^{\infty} c_{m} \int_{-\tau_{1}}^{\tau_{2}} e^{i(m-k)\omega t} = c_{m} \tau_{r} \delta_{m,k}$$

Fourier Series



$$C_{m} = \frac{1}{\tau_{r}} \sum_{\tau_{r}/2}^{\tau_{r}/2} v(t) e^{-jm\omega_{r}t} dt$$

Note that since v (t) is real

Let

$$\omega_m = m \omega_r$$

$$\sqrt[N]{(\omega_m)} = T_r C_m$$

Units of Volts/ Hz

Fourier Transforms



$$V(t) = \sum_{m=-\infty}^{\infty} \frac{1}{T_r} \tilde{V}(\omega_m) e^{j\omega_m t}$$

The spacing between adjacent frequencies is:

$$\Delta \omega_m = m \frac{2\pi}{\tau_r} - (m-1) \frac{2\pi}{\tau_r} = \frac{2\pi}{\tau_r}$$

The Fourier series becomes:

$$\mathcal{U}(t) = \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} \sqrt[N]{(\omega_m)} e^{j\omega_m t} \Delta \omega_m$$

Fourier Transforms

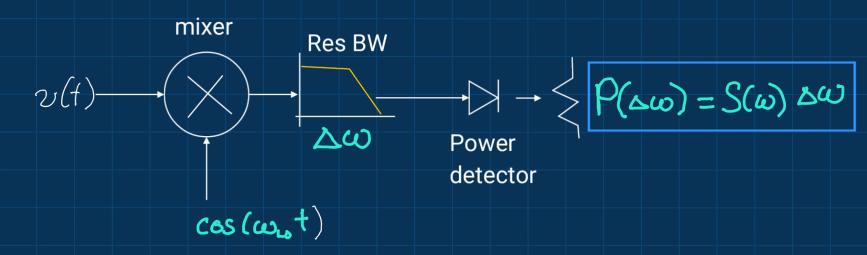


For non-periodic signals, take the limit as $T_{r} \sim (\Delta w_{r} \rightarrow 0)$

$$\nabla f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} V(\omega) e^{j\omega t} d\omega$$

$$V(\omega) = \int_{-\infty}^{\infty} v(t) e^{-j\omega t} dt$$

Swept spectrum analyzers do not measure voltages and currents. They measure power deposited into a filter of width given by the resolution bandwidth.



MAXIV RF



The power spectral density is defined as:

$$\langle \rho(t) \rangle = \frac{1}{2\pi} \int_{-\infty}^{\infty} S(\omega) d\omega$$

But time averaged power is:



Use the fact that because v(t) is real:

$$\mathcal{V}(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{V}(\omega_{i}) e^{j\omega_{i}t} d\omega_{i}$$
and
$$\mathcal{V}(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{V}(\omega_{i}) e^{j\omega_{i}t} d\omega_{i}$$

$$\delta(t-r) = \frac{1}{2\pi} \int_{\infty} e^{j\omega} (t-r) d\omega \qquad \delta(\omega-\omega') = \frac{1}{2\pi} \int_{\infty} e^{j(\omega-\omega')+} dt$$

$$\delta(\omega-\omega') = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{j(\omega-\omega')+} dt$$

results

$$\langle \rho(t) \rangle = \lim_{T \to \infty} \frac{1}{T} \int_{-\infty}^{\infty} \frac{1 V(\omega) l^2}{R} d\omega$$



The power spectral density reduces to:

$$S_f(\omega) = \lim_{T \to \infty} \frac{1}{T} |\tilde{U}(\omega)|^2$$

Since $\mathring{V}(\omega)$ has units of Volts/Hz,

$$\Sigma_{\mathbf{f}}(\omega)$$
 has units of Watts/ Hz



Consider the periodic signal again

$$V_{\rho}(t) = \sum_{n=-\infty}^{\infty} V_{r}(t-nT_{r})$$

Written as a Fourier Series
$$V_{\mu}(t) = \sum_{n=-\infty}^{\infty} C_n e^{\sqrt{n} \omega_n t}$$

Taking the Fourier transform
$$(\sqrt[n]{\omega}) = \sum_{n=-\infty}^{\infty} (n \sum_{n=-\infty}^{\infty} e^{j(n\omega_n - \omega)})^n dy$$

Results with:

$$\tilde{V}_{\rho}(\omega) = 2\pi \sum_{n=0}^{\infty} C_n \delta(\omega_{-n}\omega_r)$$



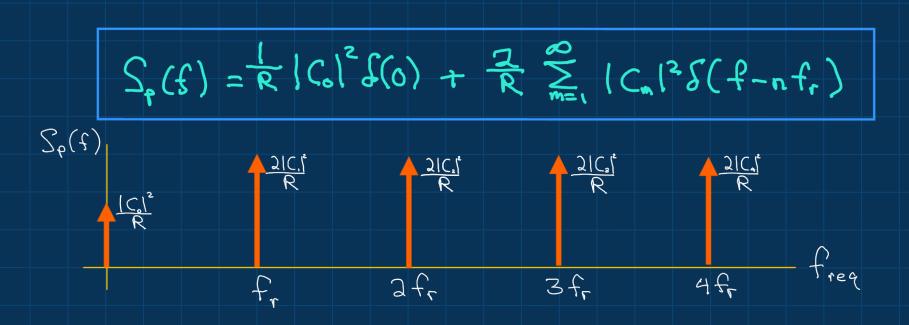
Using the magic of delta functions with the limit as T goes to infinity, the gifted reader can show:

$$S_{p}(\omega) = \frac{2\pi}{R} \sum_{m=-\infty}^{\infty} |C_{m}|^{2} \delta(\omega - m\omega_{r})$$

Since



Since spectrum analyzers do not measure phase, they do not display negative frequencies



- Note that a spectrum analyzer does not measure the power spectral density.
- It measures the power deposited in a filter of width called the resolution bandwidth at a frequency f.
- So for a periodic signal, as the resolution bandwidth is changed, the peak signal on a spectrum analyzer does not vary.

