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Laplace Transforms

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Fourier Transform Pair

$$\tilde{F}(\omega) = \int_{-\infty}^{\infty} f(t) e^{-j\omega t} dt$$

$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{F}(\omega) e^{j\omega t} d\omega$$

$\omega$  is a real number.

The problem with Fourier transforms is that they cannot handle initial conditions due to the fact  $\int_{-\infty}^{\infty}$

To handle initial conditions, consider the Laplace transform

$$\mathcal{L}[f(t)] = \tilde{f}(s) = \int_0^{\infty} f(t) e^{-st} dt$$

$s$  is a complex number!

$$s = \sigma + j\omega$$

Since  $s$  is complex number, the inverse transform is hard to compute directly



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$$\mathcal{L}^{-1}[\tilde{f}(s)] = f(t)$$

Common transform pairs

<u><math>f(t)</math></u>	<u><math>\tilde{f}(s)</math></u>
$\delta(t)$	1
$u(t)$ (step)	$1/s$
$t$	$1/s^2$
$e^{-\alpha t}$	$\frac{1}{s+\alpha}$
$\sin \omega t$	$\frac{\omega}{s^2 + \omega^2}$
$\cos \omega t$	$\frac{s}{s^2 + \omega^2}$
$e^{-\alpha t} \sin \omega t$	$\frac{\omega}{(s+\alpha)^2 + \omega^2}$
$e^{-\alpha t} \cos \omega t$	$\frac{s+\alpha}{(s+\alpha)^2 + \omega^2}$



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$$f(t)$$

$$\tilde{f}(s)$$

$$\frac{df}{dt}$$

$$s\tilde{f}(s) - f(0)$$

$$\frac{d^2f}{dt^2}$$

$$s^2\tilde{f}(s) - sf(0) - \left. \frac{df}{dt} \right|_{t=0}$$

$$\int_0^+ f(t) dt$$

$$\frac{\tilde{f}(s)}{s} + \frac{1}{s} \left[ \int_0^+ f(t) dt \right] \Big|_{t=0}$$

$$f(t-\tau)u(t-\tau)$$

$$e^{-s\tau} \tilde{f}(s)$$

### Final Value Theorem

$$\lim_{t \rightarrow 0^+} f(t) = f(0^+) = \lim_{s \rightarrow \infty} s f(s)$$

$$f(t) = e^{-\alpha t}$$

$$\tilde{f}(s) = \frac{1}{s + \alpha}$$

$$\lim_{s \rightarrow \infty} s f(s) = \lim_{s \rightarrow \infty} \frac{s}{s + \alpha} = 1$$



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Final value Theorem

$$\lim_{t \rightarrow \infty} f(t) = \lim_{s \rightarrow 0} sf(s)$$

(only if  $f(t)$  is stable)

Uncharged Capacitor

$$i(t) = C \frac{dv}{dt}$$

$$\tilde{i}(s) = sC \tilde{v}(s) - C v(0)$$

Uncharged Inductor.

$$v(t) = L \frac{di}{dt}$$

$$\tilde{v}(s) = sL \tilde{i}(s)$$



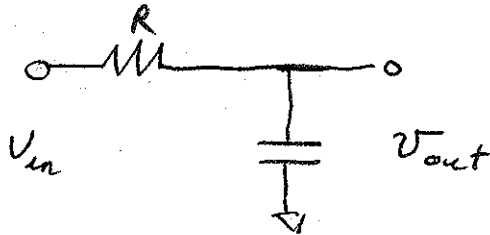
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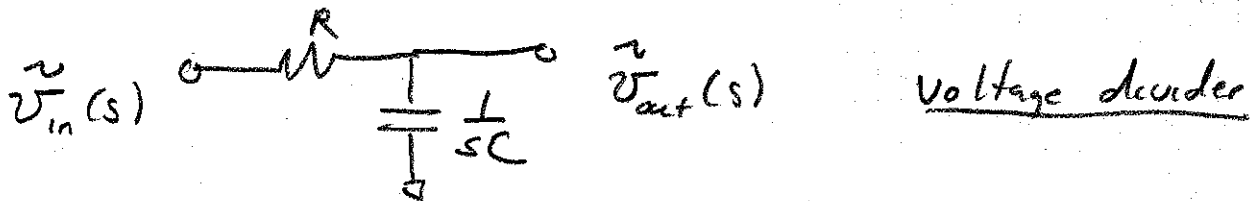
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## Low Pass Filter



In Laplace Domain (with everything uncharged)



$$V_{out}(s) = \frac{\frac{L}{sC}}{R + \frac{L}{sC}} V_{in}(s)$$
$$= \frac{1}{1 + sRC}$$

What is the frequency response

$$s \Rightarrow j\omega$$

$$V_{out}(j\omega) = \frac{1}{1 + j\omega RC} V_{in}(j\omega)$$



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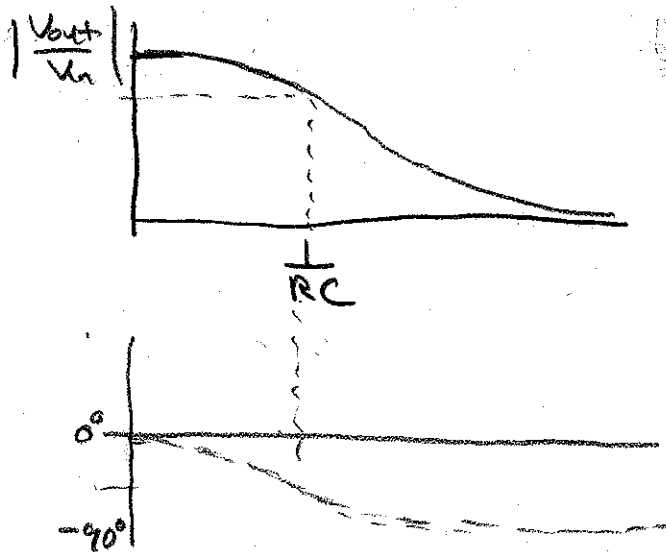
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When  $\omega = \frac{1}{RC}$

$$\frac{V_{out}}{V_{in}} = \frac{1}{1+j} = \frac{1-j}{2}$$

$$\left| \frac{V_{out}}{V_{in}} \right| = \frac{1}{\sqrt{2}}$$

$$\angle \left( \frac{V_{out}}{V_{in}} \right) = -45^\circ$$





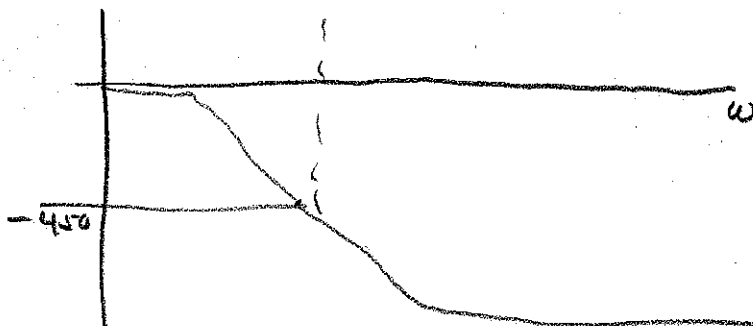
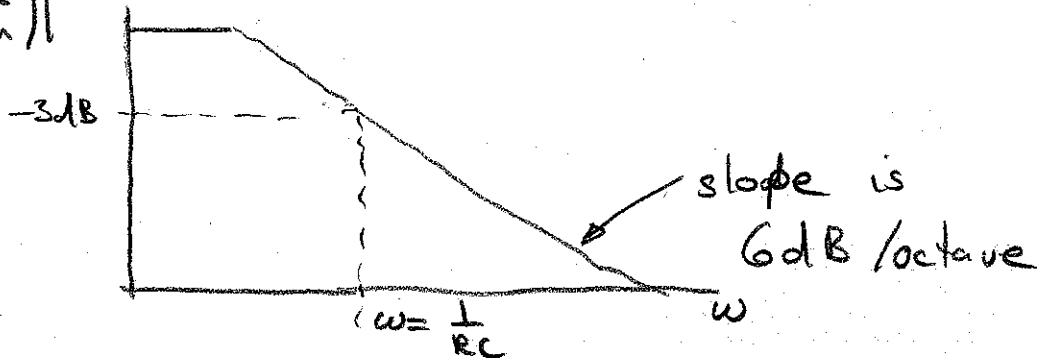
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$$10 \log_{10} \left( \left| \frac{V_{out}}{V_{in}} \right|^2 \right)$$



The 3dB bandwidth

$$\Delta\omega = \frac{1}{RC}$$

$$\Delta f = \frac{1}{2\pi RC}$$

What does bandwidth mean?



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$$\tilde{V}_{out}(s) = \frac{1}{1+sRC} \tilde{V}_{in}(s)$$

Let  $V_{in}(t) = V_0 \Delta T \delta(t)$  (get units right)  
1/time

$$\tilde{V}_{in}(s) = V_0 \Delta T$$

$$\tilde{V}_{out}(s) = \frac{V_0 \Delta T}{1+sRC}$$

$$\tilde{V}_{out}(s) = \frac{V_0 \Delta T}{RC} \left( \frac{1}{s + \frac{1}{RC}} \right)$$

Let  $\Delta T = RC$

$$\tilde{V}_{out}(s) = V_0 \left( \frac{1}{s + \frac{1}{\Delta T}} \right)$$

$$V_{out}(t) = V_0 e^{-t/\Delta T}$$





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Now consider a second impulse on the circuit.

$$v_{in_2}(t) = v_0 \Delta T \delta(t - \tau_d) u(t - \tau_d)$$

$$v_{in_2}(s) = v_0 \Delta T e^{-s \tau_d}$$

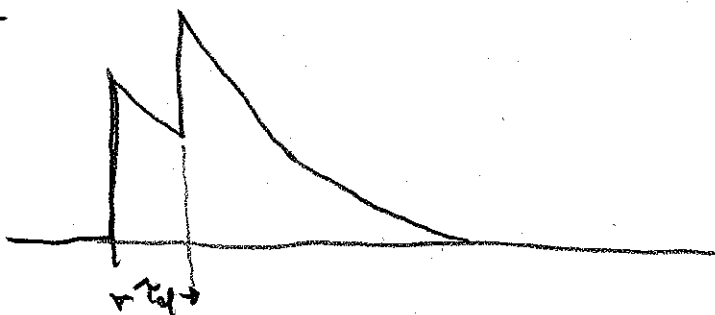
$$v_{out_2}(s) = v_0 \frac{e^{-s \tau_d}}{(s + 1/\Delta T)}$$

$$v_{out_2}(t) = v_0 e^{-(t - \tau_d)/\Delta T} u(t - \tau_d)$$

Now consider the two impulses together

$$v_{in_{1+2}}(t) = v_0 \Delta T \delta(t) + v_0 \Delta T \delta(t - \tau_d) u(t - \tau_d)$$

$$v_{out_{1+2}}(t) = v_0 e^{-t/\Delta T} + v_0 e^{-(t - \tau_d)/\Delta T} u(t - \tau_d)$$





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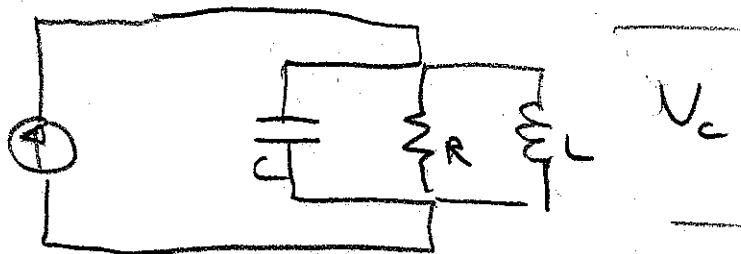
If  $\tau_d \ll \Delta T$  we cannot distinguish the two impulses on the output

$\Delta T$  is a measure of our ability to resolve the two impulses

$$\Delta T = \frac{1}{2\pi \Delta f}$$

∴ Bandwidth is a measure of how well a circuit can resolve impulses

Consider the cavity circuit



$$Z(s) = \frac{1}{sC + \frac{1}{R} + \frac{1}{sL}}$$



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$$Z(s) = \frac{s/L}{s^2 + \frac{s}{RC} + \frac{1}{LC}}$$

Define  $\omega_0^2 = \frac{1}{LC}$

$$Q = \omega_0 RC$$

$$Z(s) = \frac{s \omega_0 R/Q}{s^2 + s \frac{\omega_0}{Q} + \omega_0^2}$$

$$I_{\text{gen}}(t) = \int_{-\infty}^t q \delta(t) dt$$

$\rho_L$   
charge

Remember

$$\mathcal{L}(e^{-\alpha t} \sin \omega_r t) = \frac{\omega_r}{(s+\alpha)^2 + \omega_r^2}$$

$$\mathcal{L}(e^{-\alpha t} \cos \omega_r t) = \frac{s+\alpha}{(s+\alpha)^2 + \omega_r^2}$$



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$$\mathcal{L} \left( e^{-\alpha t} \cos \omega_r t - \frac{\alpha}{\omega_r} e^{-\alpha t} \sin \omega_r t \right) = \frac{s}{(s+\alpha)^2 + \omega_r^2}$$

Comparing -

$$2\alpha = \frac{\omega_0}{Q}$$

$$\alpha = \frac{1}{2} \frac{\omega_0}{Q}$$

$$\omega_0^2 = \alpha^2 + \omega_r^2$$

$$\omega_r^2 = \omega_0^2 \left( 1 - \frac{1}{4Q^2} \right)$$

$$v(t) = q \omega_0 \frac{R}{Q} e^{-\frac{\omega_0}{2Q} t}$$

$$\cdot \left[ \cos \omega_r t - \frac{1}{\sqrt{4Q^2 - 1}} \sin \omega_r t \right]$$

Note

- 1) Cavity rings down at  $\tau = \frac{2Q}{\omega_0}$
- 2) Ring frequency slightly shifted down from resonant freq.
- 3) Amplitude  $\propto R/Q$  not  $R$ !



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Look at

$$\mathcal{L}(e^{-\alpha t} \sin \omega_r t) = \frac{\omega_r}{(s + \alpha)^2 + \omega_r^2}$$

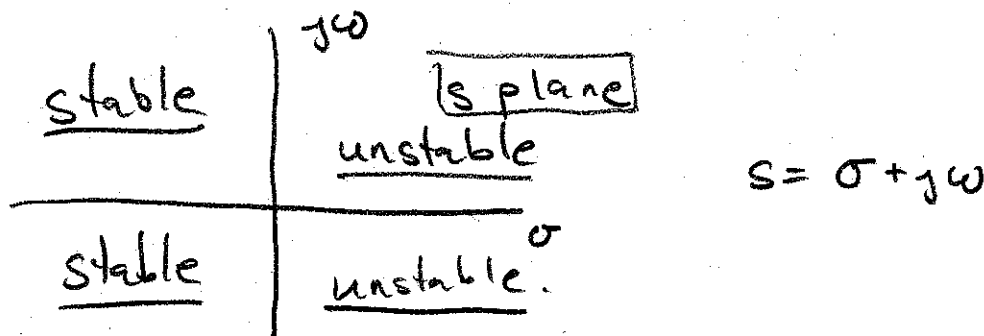
$$(s + \alpha)^2 + \omega_r^2 = (s - (j\omega_r - \alpha))(s - (-j\omega_r - \alpha))$$

$\nearrow$   
 poles of the response

(places where denominator goes to zero)

The characteristic frequencies of the system is determined by the poles of the Laplace Transform.

Note that for the system to be stable, the real part of the poles must be  $< 0$ . (i.e., the poles must be on the left hand side of the s-plane)







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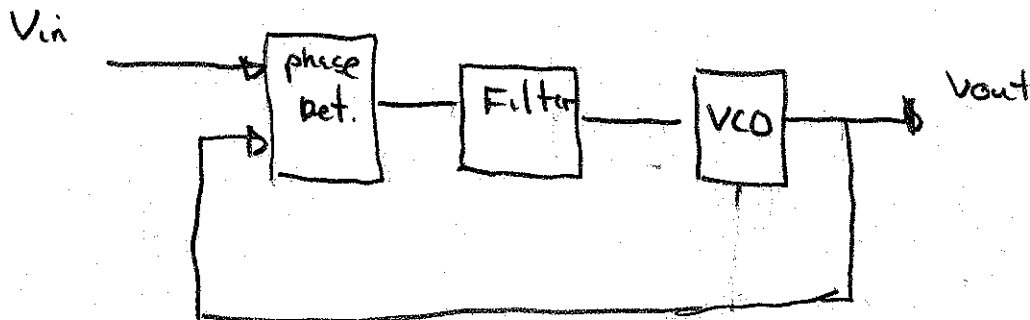
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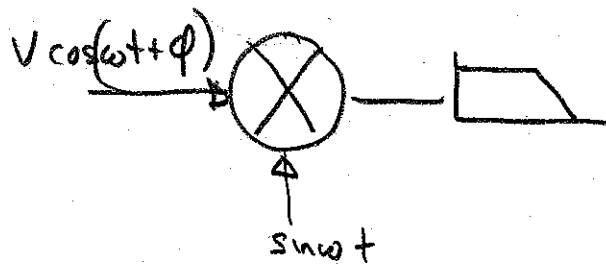
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if  $G > 1$  system will be unstable!

## Phase Lock Loop



How do you make a phase detector?

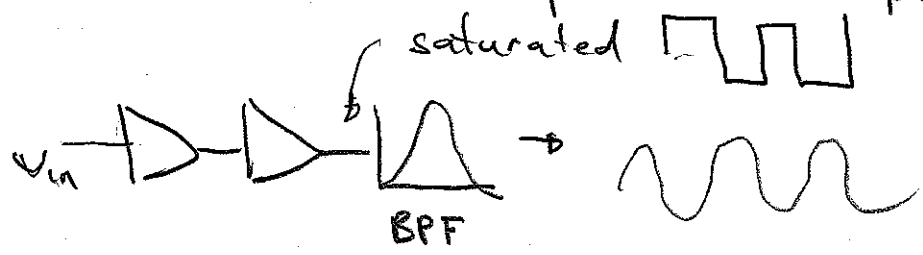


$$V_{\text{mixer}} = V \cos(\omega t + \phi) \sin \omega t$$

$$= \frac{1}{2} V \sin(2\omega t + \phi) + \frac{V}{2} \sin(\phi)$$

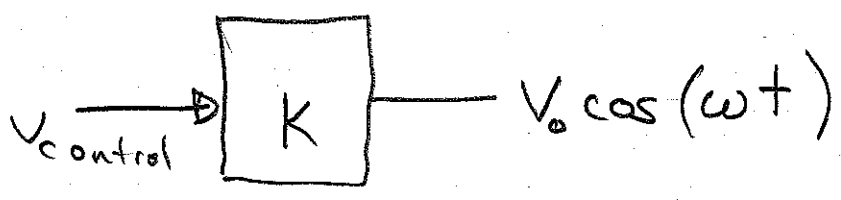
$$V_{\text{filter}} = \frac{V}{2} \sin \phi$$

To remove Amplitude dependence



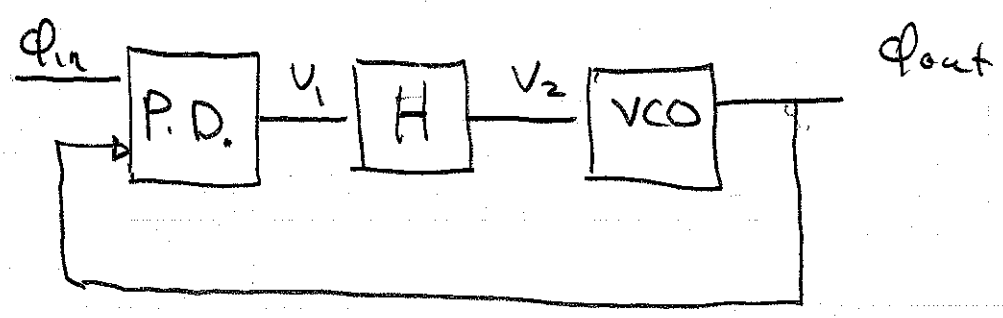
A VCO is a voltage controlled oscillator

The frequency out of the oscillator is proportional to the control voltage



$$\omega = 2\pi K V_c$$

Lets use the phase as our variable



$$V_1 = \phi_{in} - \phi_{out}$$

$$V_2 = H(s) V_1$$



$$\phi_{out} = \int 2\pi f_{out} dt$$

$$f_{out} = K V_2$$

$$\phi_{out} = \frac{2\pi K}{s} V_2$$

$$\phi_{out} = \frac{2\pi K}{s} H(s) (\phi_{in} - \phi_{out})$$

$$= 2\pi K \frac{H}{s} \phi_{in} - 2\pi K \frac{H}{s} \phi_{out}$$

$$\phi_{out} = \frac{2\pi K \frac{H}{s}}{1 + 2\pi K \frac{H}{s}} \phi_{in}$$

$$\text{Let } H(s) = \frac{1}{1 + s\tau_h}$$

$$G = 2\pi K$$

$$\phi_o = \frac{G}{s(1+s\tau_h) \left( 1 + \frac{G}{s(1+s\tau_h)} \right)}$$

$$= \frac{G}{s^2\tau_h + s + G}$$

$$\phi_0 = \frac{G/\tau_h}{s^2 + \frac{s}{\tau_h} + \frac{G}{\tau_h}} \quad \text{lin}$$

$$s_p = -\frac{1}{2\tau_h} \pm \left( \left( \frac{1}{2\tau_h} \right)^2 - \frac{G}{\tau_h} \right)^{1/2}$$

Critical damping

$$G = \frac{1}{4\tau_h^2}$$

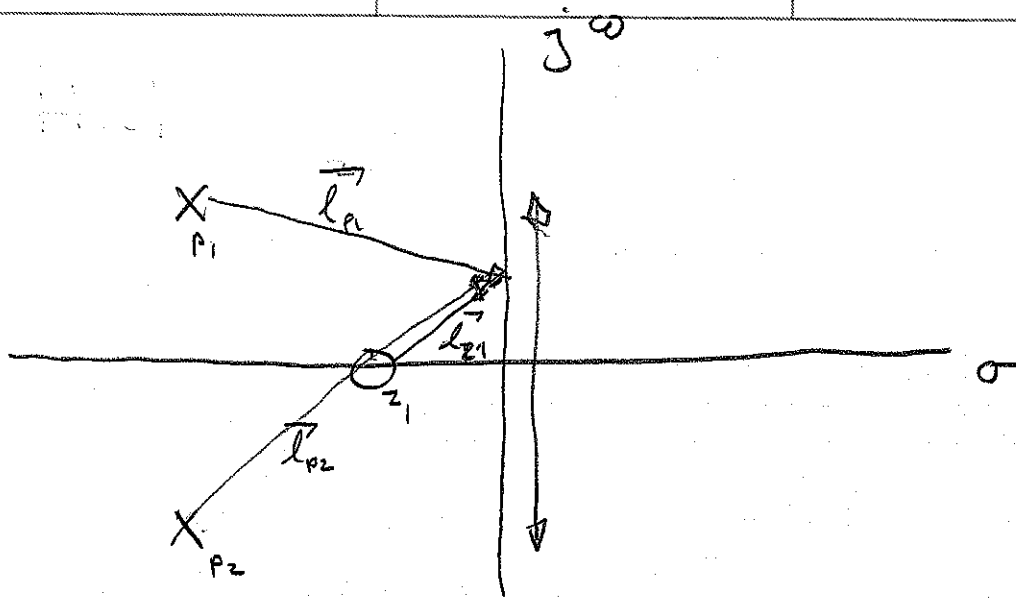
$$K = \frac{1}{8\tau_h} \frac{K}{\tau_h}$$

$$\phi_0 = \frac{\left( \frac{1}{2\tau_h} \right)^2}{\left( s + \frac{1}{2\tau_h} \right)^2} \quad \text{lin}$$

Plotting the frequency response

In General

$$H(s) = \frac{(s-s_{z_1})(s-s_{z_2}) \dots (s-s_{z_{n-1}})}{(s-s_{p_1})(s-s_{p_2}) \dots (s-s_{p_n})}$$

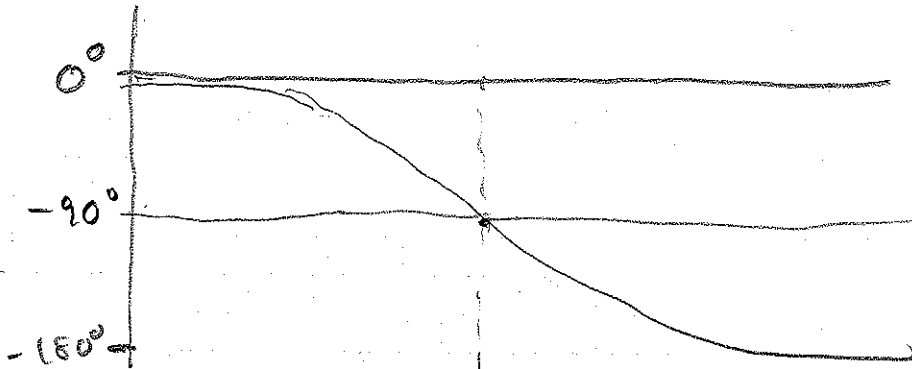
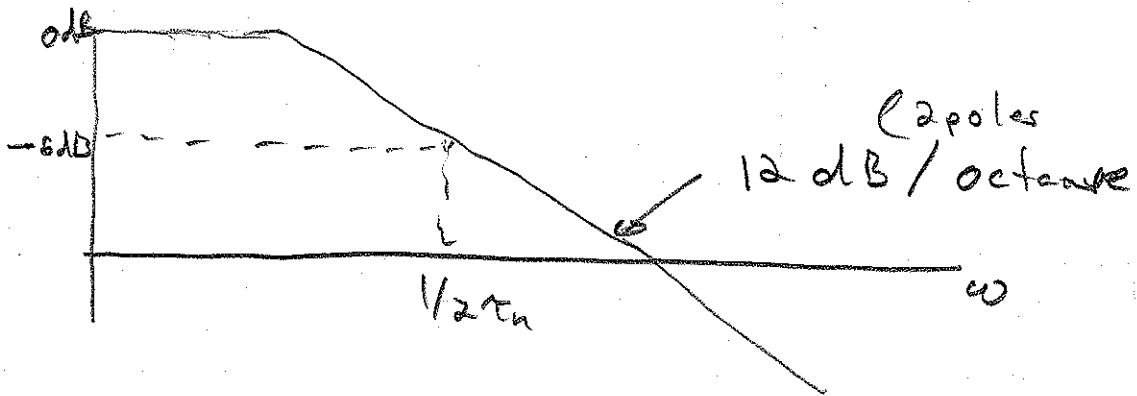
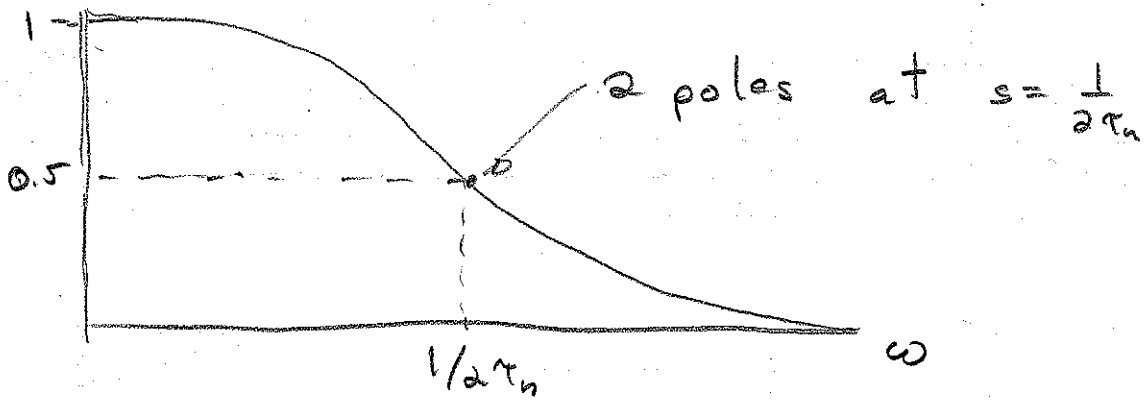
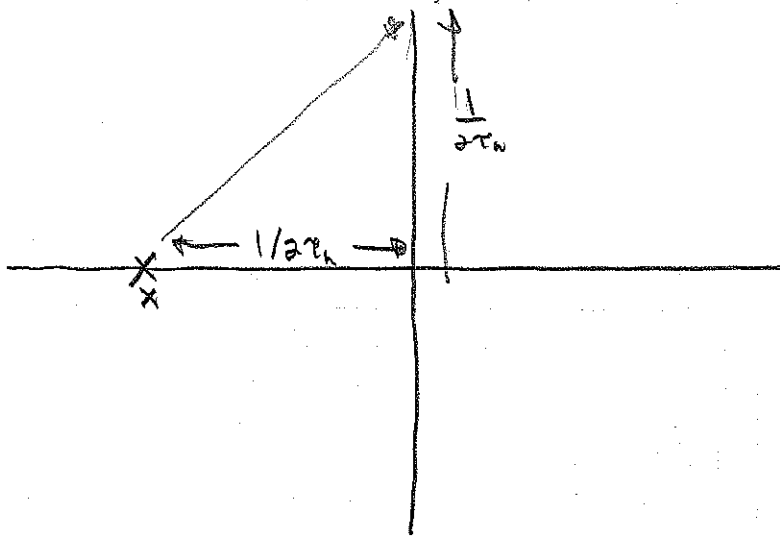


$$|H(j\omega)| = \frac{\prod |\vec{l}_{z_i}|}{\prod |\vec{l}_{p_i}|}$$

$$\angle H(j\omega) = \sum \angle \vec{l}_{z_i} - \sum \angle \vec{l}_{p_i}$$

Look at phase lock loop response

$$\phi_{out}(s) = \frac{\left(\frac{1}{2\tau_n}\right)^2}{\left(s + \frac{1}{2\tau_n}\right)^2} \phi_{in}$$





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$$V_2 = \frac{s}{G} \phi_{out}$$

For critical damping

$$G = \frac{1}{4\tau_n}$$

$$V_2 = s \cdot 4\tau_n \phi_{out}$$

$$V_2 = \frac{s \left(\frac{1}{2\tau_n}\right)^2 4\tau_n \phi_{in}}{\left(s + \frac{1}{2\tau_n}\right)^2}$$

$$= \frac{s / \tau_n}{\left(s + \frac{1}{2\tau_n}\right)^2} \phi_{in}$$

